## Week (5

PHY 106 Quantum Physics
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## Book: Wave "groups"

2.3) Wave packets and dispersion

Movies: Elm waves are also particles
What if traditional particles
(electrons) are also waves?
But particle can be in a specific place
Extended travelling wave? Is not!!!
https://phet.colorado.edu/sims/html/wave-on-a-string/latest/wave-on-a-string_en.html

## But in our rope app/experiment, can also see wave pulse:



## How to get from sin/cos wave to pulse?



It is useful to keep discussing (also pulse) waves in terms of sin and cos, since all our week 3 material will apply!

Idea: use superposition principle to combine different sine waves into hump?
2.3.1.) Beating of two waves

Recall question from tutorial 2 :
Adding two sines with slight wavelength difference Middle bit-> pulse


## Beating of two waves



## Beating of two waves

Let's do the math:


$$
y_{t o t}(x)=y_{1}(x)+y_{2}(x)
$$

$$
y_{1}(x)=A \cos \left(\frac{2 \pi}{\lambda_{1}} x\right)+y_{2}(x)=A \cos \left(\frac{2 \pi}{\lambda_{2}} x\right)
$$

Use
$\cos (a)+\cos (b)=2 \cos \left(\frac{a+b}{2}\right) \cos \left(\frac{a-b}{2}\right)$

$$
y_{\text {tot }}(x)=2 A \cos \left[\left(\frac{\pi}{\lambda_{1}}+\frac{\pi}{\lambda_{2}}\right) x\right] \cos \left[\left(\frac{\pi}{\lambda_{1}}-\frac{\pi}{\lambda_{2}}\right) x\right]
$$

## Beating of two waves

$y_{\text {tot }}(x)=2 A \cos \left[\left(\frac{\pi}{\lambda_{1}}-\frac{\pi}{\lambda_{2}}\right) x\right] \cos \left[\left(\frac{\pi}{\lambda_{1}}+\frac{\pi}{\lambda_{2}}\right) x\right]$

## Envelope:

$k_{\text {low }}=\left(k_{1}-k_{2}\right) / 2$
Carrier:


## Beating of two waves

$\begin{aligned} & \text { Total: } \\ & y_{\text {tot }}(x)=2 A \cos \left[\left(\frac{\pi}{\lambda_{1}}-\frac{\pi}{\lambda_{2}}\right) x\right] \cos \left[\left(\frac{\pi}{\lambda_{1}}+\frac{\pi}{\lambda_{2}}\right) x\right]\end{aligned}$
Envelope:

$$
k_{\text {low }}=\left(k_{1}-k_{2}\right) / 2 \quad k_{\text {high }}=\left(k_{1}+k_{2}\right) / 2
$$

Carrier:

## Beating of two waves

Total:
$\underline{y_{t o t}(x)}=2 A \cos \left[\left(\frac{\pi}{\lambda_{1}}-\frac{\pi}{\lambda_{2}}\right) x\right] \cos \left[\left(\frac{\pi}{\lambda_{1}}+\frac{\pi}{\lambda_{2}}\right) x\right]$

## Envelope:

$k_{\text {low }}=\left(k_{1}-k_{2}\right) / 2$

Carrier:
$k_{\text {high }}=\left(k_{1}+k_{2}\right) / 2$

Note: The math above simply gives us a product of two cosines. Which we call carrier and which envelope, simply is decided by which one has the much larger wavelength (that cosine is called envelope)

### 2.3.2) Fourier decomposition

We managed to make wave more "pulsey" by adding two.
We can perfectly form any function if we take more waves:

## Fourier theorem:

Any even function $f(x)$ can be written as:

$$
\begin{equation*}
f(x)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} d k \tilde{g}(k) \cos (k x) \tag{42}
\end{equation*}
$$

If $\mathrm{f}(\mathrm{x})$ is periodic with period L

$$
\begin{equation*}
f(x)=\sum_{n=0}^{\infty} g_{n} \cos \left(\frac{2 \pi n}{L} x\right) \tag{43}
\end{equation*}
$$

## Fourier decomposition

## Fourier theorem:

Any even function $f(x)$ can be written as:

$$
\begin{equation*}
f(x)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} d k \tilde{g}(k) \cos (k x) \tag{42}
\end{equation*}
$$

If $f(x)$ is periodic:

$$
\begin{equation*}
f(x)=\sum_{n=0}^{\infty} g_{n} \cos \left(\frac{2 \pi n}{L} x\right) \tag{43}
\end{equation*}
$$

The coefficients $g_{n}$ can be found via:

$$
\begin{equation*}
g_{n}=\frac{2}{L} \int_{-L / 2}^{L / 2} d x f(x) \cos \left(\frac{2 \pi n}{L} x\right) \tag{44}
\end{equation*}
$$

## BONUS MATERIAL: Fourier decomposition

Of course it also works for odd functions $f(x)=-f(-x)$, using sines

$$
\begin{align*}
& f(x)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} d k \tilde{g}(k) \sin (k x)  \tag{45}\\
& f(x)=\sum_{n=0}^{\infty} h_{n} \sin \left(\frac{2 \pi n}{L} x\right) \quad h_{n}=\frac{2}{L} \int_{-L /}^{L / 2} d x f(x) \sin \left(\frac{2 \pi n}{L} x\right)
\end{align*}
$$

## BONUS: Fourier decomposition

More generally any function $f(x)$ can be written as:

$$
\begin{equation*}
f(x)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} d k \tilde{f}(k) e^{i k x} \tag{48}
\end{equation*}
$$

with

$$
\begin{equation*}
\tilde{f}(k)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} d x f(x) e^{-i k x} \tag{49}
\end{equation*}
$$

using

$$
\begin{align*}
& e^{i k x} \equiv \cos (k x)+ \underbrace{i} \sin (k x)  \tag{50}\\
&=\sqrt{-1}
\end{align*}
$$

## Fourier decomposition

Example: Sawtooth curve (see fourie__deoomposition_v1.m)


Discard this ( $g_{n}$ small), since wavelength wrong See webapps in tutorial 5, for this!!!

## Fourier decomposition

## Example: Sawtooth curve



Keep this ( $\mathrm{g}_{\mathrm{n}}$ large), since wavelength good

## Fourier decomposition

## Example: Sawtooth curve



Need shorter wavelengthes for small features

To see the complete animation, follow this link: http://home.iiserb.ac.in/~sebastian/ material/QuantPhys/fourier sawtooth.mp4

Legend: (see Eq. 43 and 44)
black line $\quad f(x)=$ sawtooth function
red dashed
cumulative sum $s_{n}(x)=\sum_{m=0}^{n} g_{m} \cos \left(\frac{2 \pi m}{L} x\right)$
blue or magenta
$g_{n}<10^{-5} \quad g_{n}>10^{-5}$
trial cosine $\cos \left(\frac{2 \pi n}{L} x\right)$

### 2.3.3) Gaussian wave packet

We call the combination of many waves a wave packet

A neat case is the Gaussian wave packet

## Math: Gaussian function

$$
g(x)=\frac{1}{\left.\sqrt{2 \pi}\right|_{x}} e^{-\frac{\left(x-x_{0}\right)^{2}}{\left|2 \sigma_{x}\right|}}
$$



- $\sigma_{x}$ is called the width (or standard deviation) of the Gaussian
- Pre-factor fixes normalisation

$$
\int_{-\infty}^{\infty} d x g(x)=1
$$

## Gaussian wave packet

We call the combination of many waves a wave packet

A neat case is the
Gaussian wave packet

$$
\begin{align*}
& f(x)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} d k \tilde{g}(k) \cos (k x) \quad \text { repeat (42) } \\
& \tilde{g}(k)=\frac{1}{\sqrt{2 \pi} \sigma_{k}} e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \tag{51}
\end{align*}
$$

Gaussian wave packet

$$
\begin{aligned}
& f(x) \sim \int_{-\infty}^{\infty} d k e^{-\frac{\left(\alpha-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \cos (k x) \\
& f(x) \sim e^{-\frac{x^{2}}{2 \sigma_{x}^{2}}} \cos \left(k_{0} x\right)
\end{aligned}
$$

Envelope of Gaussian-w.p. is again Gaussian:


To see the complete animation, follow this link: http://home.iiserb.ac.in/~sebastian/ material/QuantPhys/fourier gaussian carrier.mp4

Legend: (see Eq. 43 and 44) black line

$$
f(x) \sim e^{-\frac{x^{2}}{2 \sigma_{x}^{2}}} \cos \left(k_{0} x\right)
$$

red dashed
cumulative sum $s_{n}(x)=\sum_{m=0}^{n} g_{m} \cos \left(\frac{2 \pi m}{L} x\right)$
blue or magenta
$g_{n}<10^{-5} \quad g_{n}>10^{-5} \quad$ trial cosine $\cos \left(\frac{2 \pi n}{L} x\right)$

Gaussian wave packet (math version)
factor to give overall scale of function

$$
f(x)=\mathscr{N} \int_{-\infty}^{\infty} d k e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \cos (k x)
$$

We can solve the above integral and find:

$$
f(x)=\mathcal{N} \sqrt{2 \pi} \sigma_{k} e^{-\frac{x^{2} \sigma_{k}^{2}}{2}} \cos \left(k_{0} x\right)
$$

We can write this as

$$
f(x)=\tilde{\mathscr{N}} e^{-\frac{x^{2}}{2 \sigma_{x}^{2}}} \cos \left(k_{0} x\right) \quad \text { with } \quad \sigma_{x}=1 / \sigma_{k}
$$

...as we did on the previous slide.

Gaussian wave packet

$$
f(x) \sim \int_{-\infty}^{\infty} d k e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \cos (k x)
$$

Envelope of Gaussian-w.p. is again Gaussian:


Gaussian wave packet

$$
f(x) \sim \int_{-\infty}^{\infty} d k e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \cos (k x)
$$




Gaussian: $x$ and $k$ widths are inverse

$$
\begin{equation*}
\sigma_{x}=1 / \sigma_{k} \tag{52}
\end{equation*}
$$

- If we want a more localised wave packet, we need a larger range of wave lengths!!!


## Examples:



### 2.3.4) Dispersion

Dispersion For waves in a medium, the phase velocity V may depend on the wave frequency $\omega$.

- In other words, the relation between $\omega$ and k is not proportional (as in $\omega=V k$ )
-Then phase velocity $V=\omega / k$ is not constant
We call the dependence of $\omega$ on $k$ Dispersion relation $\omega=f(k)$


## Dispersion

-Wave Eqn. (13) predicts equal phase velocity $\mathrm{V}=$ const. for all waves

$$
\frac{\partial^{2}}{\partial x^{2}} y(x, t)=\frac{1}{V^{2}} \frac{\partial^{2}}{\partial t^{2}} y(x, t)
$$

-Thus if we have dispersion it needs modification, e.g.

$$
\frac{\partial^{2}}{\partial x^{2}} y(x, t)-\alpha \frac{\partial^{4}}{\partial x^{4}} y(x, t)=\frac{1}{\beta^{2}} \frac{\partial^{2}}{\partial t^{2}} y(x, t)
$$

(I don't tell you what $\alpha, \beta$ are, this is just an example for the mathematical structure)

## Dispersion


https://www.youtube.com/watch?v=KbmOcT5sX7I

## Dispersion

## Example:

-The dependence of phase velocity on frequency/ wavenumber is often weak i.e. $V(\omega) \approx V_{0} \quad \forall \omega$
C/V Lanthanum dense flint LaSF9

### 2.3.5) Group velocity

Consider Gaussian wave-packet

$$
f(x) \sim \int_{-\infty}^{\infty} d k e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \cos (k x)
$$

Now lets make waves moving

$$
f(x, t) \sim \int_{-\infty}^{\infty} d k e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}}} \frac{\cos (k x-\omega t)}{\cos [k(x-V t)]}
$$

If no dispersion: $\omega=V k$ (see Eq. 8, same V)

## Group velocity

Moving Gaussian wave-packet without dispersion


### 2.3.5) Group velocity

Consider Gaussian wave-packet

$$
f(x) \sim \int_{-\infty}^{\infty} d k e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \cos (k x)
$$

Now lets make waves moving

$$
f(x, t) \sim \int_{-\infty}^{\infty} d k e^{-\frac{\left(k-k_{0}\right)^{2}}{2 \sigma_{k}^{2}}} \cos [k(x-\underline{V(k) t)}]
$$

If dispersion: Different velocities $V(k)$ for different " $k$ " parts of wavepacket

## Group velocity

## Moving Gaussian wave-packet


see more animations here:
https://blog.soton.ac.uk/soundwaves/further-concepts/2-dispersive-waves/

## Simpler example: Motion of beating waves

Go back to beating example

Assume two
 different waves

$$
\begin{aligned}
& y_{1}(x, t)=A \cos [(\omega-\Delta \omega / 2) t-(k-\Delta k / 2) x] \\
& y_{2}(x, t)=A \cos [(\omega+\Delta \omega / 2) t-(k+\Delta k / 2) x]
\end{aligned}
$$

second has only slightly different $\omega$ and k .

## Motion of beating waves

Add them as in section 2.3.1.)

$$
\begin{align*}
y(x, t) & =y_{1}(x, t)+y_{2}(x, t) \\
& =2 A \cos \left(\frac{\Delta \omega}{2} t-\frac{\Delta k}{2} x\right) \cos (\omega t-k x) \tag{53}
\end{align*}
$$

Moving Envelope: Moving Carrier:

- See book for details.


## Movie: Motion of beating waves



Motion of beating waves
$y=\frac{2 A \cos \left(\frac{\Delta \omega}{2} t-\frac{\Delta k}{2} x\right)}{\text { Moving Envelope: }} \frac{\cos (\omega t-k x)}{\text { Moving Carrier: }}$
Using Eq. (8) we infer for motion of envelope

$$
\begin{array}{ll}
\text { Group velocity } \\
\qquad \begin{aligned}
v_{g}=\frac{\Delta \omega}{\Delta k} & \text { (Two waves) } \\
v_{g}=\frac{d \omega}{d k} & \text { (Many waves }
\end{aligned}
\end{array}
$$

Movie: Gaussian moving wavepacket


